Nanolithography based on an atom pinhole camera for fabrication of metamaterials

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Abstract

We have experimentally realized a method of images construction in atom optics, based on the idea of optical pinhole camera. Generation of identical images with maximum resolution has been explored. With the use of an atom pinhole camera we have built on a Si and glass surfaces an array of identical arbitrary-shape atomic nanostructures with the minimum size of an individual nanostructure’s element down to 50 nm. Limitations of the approach for fabrication of metamaterials are discovered.

Keywords: Atom nanolithography; Atom optics; Metamaterials

1. Introduction

There is a rising interest in design and discovery of metamaterials [1–3]. In the optical frequency range metamaterials could be made on the basis of artificially created atomic and molecular structures on the surface, with characteristic size of individual structure element in nanometer range [3]. At present, the most developed method for surface nanostructure creation is optical photolithography [4]. Photolithography, or exposure of light on a photosensitive material through a photomask, is a widespread technique used to replicate patterns. It is highly developed and well-suited for applications in microelectronics [5]. Today, photolithography makes possible nanostructures with minimum lateral dimensions down to 45 nm. It is, however, limited to photosensitive materials and is suitable only for fabrication on planar surfaces. Another problem is that in all conventional optical techniques the resolution is restricted by diffraction. When in the path of the light there is an aperture smaller than approximately one half of its wavelength $\lambda$, diffraction occurs. In the context of lithography, this means that unlimited reduction of structure size is not possible in mask-based processes: when a gap in a mask becomes comparable with $\lambda/2$, the contours of resulting structures will no longer be clearly defined because of the diffraction effect. Utilization of light sources with shorter wavelengths solves the problem, but makes the method more complicated and expensive. Besides, the light with short wavelengths imposes physical limitations on materials for optical elements (lenses, mirrors, phase masks, etc.).
Nanolithographic methods, based on the use of material-particles optics instead of light optics, enables the problem of diffraction limit to be solved, because for most of the particles de Broglie wavelength is essentially less than 1 nm. At present, nanolithography based on utilization of focused beams of charged particles (electrons or ions) is best developed [6]. Use of neutral particles instead of charged ones for nanolithography offers few side benefits. Firstly, the lack of charge removes the problem of Coulomb repulsion. Secondly, low kinetic energy of atoms allows to create nanostructures on a substrate without destruction of its surface, what in turn makes it possible to use as substrates a wider class of surfaces: biomaterials, electric microcircuits, etc. Thirdly, the utilization of neutral particles enables to realize the “direct method” of nanolithography: nanostructures are created just from the required material.

Nanolithography on the basis of neutral atoms is not so well developed as that using light or charged particles. Different approaches to nanostructure creation based on the effect of surface self-assembly of atoms [7], stencil mask nanolithography [8–10], individual atoms control on a surface through the use of a tunnel microscope [11] are known. The above-listed methods have several restrictions on material, form and linear dimensions reproduction accuracy of nanostructures to be created.

An alternative for neutral particles nanolithography is atom optics [12–15]. During past 10–15 years, atom optics has developed into an important subfield of atomic, molecular and optical physics, and contributes to different areas of technology [13]. One of the important trends in atom optics is development of basic elements, which are similar to familiar devices of conventional light optics, such as atom lenses, mirrors, beam splitters and interferometers, as well as application of these elements in practical devices. Among many possible applications of atom–optical elements, a potentially important one is micro- and nanofabrication of material structures, usually referred to as atom lithography [13].

In the method, internal and external atomic degrees of freedom are controlled with a very high precision by external electromagnetic fields (or material structures) and thus results in high-resolution surface patterning. Methods of atom lithography are founded on deposition of atoms from a beam sharply focused by an atom lens, generated by a spatially inhomogeneous field of laser radiation [16,17]. Despite numerous suggestions and experimental studies in atom beam focusing [18], the issue has not been resolved experimentally. The central problem is generation of an atom–electromagnetic field interaction potential, which in properties would be close to “ideal” lens for atoms: with minimum chromatic aberration and compensated astigmatism while permitting to focus the atom beam into a spot, diffractionally limited in space.

Recently new approach for nanostructures creation, based on the idea of object imaging in atom optics via atom pinhole camera was demonstrated (APC nanolithography) [19,20]. The approach has several important advantages: (1) it makes possible nanostructures with typical size down to 30 nm; (2) the nanostructures can have an arbitrary prearranged shape; (3) size and form of nanostructures are determined by well-controlled parameters.

This research has been the first to analyze the possibility of use APC nanolithography for creation of metamaterials.

2. Nanolithography based on atom pinhole camera

In an atom pinhole camera, atoms act as photons in an optical pinhole camera and therefore the main principles of imaging by an atom pinhole are akin to those used in light optics of a pinhole camera. As is generally known from light optics, a pinhole camera is capable of producing high-quality (distortion free and high resolution) object images. Two major questions should be answered in constructing particular pinhole camera model: (1) what is the optimum size of the pinhole to attain maximum resolution; (2) what resolution in this case is expected. From qualitative physical considerations it is obvious that, at given distance to the image plane, a large pinhole does not allow to gain an image of high quality. On the other hand, with far too small an aperture the diffusion of atoms also hinders an image construction. The standard approach to imagery through the use of pinhole camera is to consider image construction of a point object at infinity. In this case a plane wave is incident on a screen with pinhole of radius \( s \) and at distance \( l \) (focal length of the pinhole camera) a spot of radius \( r_g \) is generated. When the screen pinhole is large, the spot presents its geometrical shadow, and the image radius equals that of the pinhole. As the pinhole decreases, the image spot must be described by physical optics and Fresnel (or Fraunhofer) diffraction pattern of the pinhole. In this case, for a circular pinhole the spot radius \( r_d \approx 0.61 \lambda / s \). Hence the radius \( R \) of the image spot made by pinhole camera is roughly (in the axial approximation) the sum of the image geometrical radius \( r_g \) and the radius of the diffraction pattern caused by the aperture \( R = r_g + 0.61 \lambda (l/s) \), where \( l \) is the distance between pinhole camera and image plane. The smallest image is achieved when geometrical optics and theory
of diffraction give the same results, i.e. when the condition \( s^2 \approx 0.61 \lambda l \) is fulfilled. Closer examination based on theory of diffraction shows that the resolution of a pinhole camera can be even better than the geometrical one. Precise calculations [21] show that the image spot diameter at the optimum distance is three times smaller than the pinhole diameter.

An atom pinhole camera (like an optical pinhole camera) is free from linear distortion aberration. The lack of linear distortion follows from the argument based on Fermat’s principle (for small aperture) and from ray optics treatment (in geometrical approximation). The pinhole camera astigmatism comes about because the pinhole aperture appears as an ellipse when viewed not at right angle. The optimum focal length in one plane then differs from that in the perpendicular plane. An atom pinhole camera is also prone to chromatic aberration. This is evident from the relationship between focal distance and wavelength: \( l_{\text{opt}} \approx s^2 / \lambda dB \). In material-particles optics, for the lenses based on electromagnetic interaction potentials the relationship between chromatic aberration and velocity of particles is quadratic. In an atom pinhole camera by virtue of linear relationship between optimum focal length and velocity of an atom, chromatic aberration is linear with respect to the atom velocity, i.e. for atom pinhole cameras this type of aberration is of lesser importance.

The preceding analysis of atom pinhole camera presupposes an infinitely thin screen. In a real experiment the screen thickness is finite, and at sufficiently small aperture the action of van der Waals forces takes effect in atom’s motion through the pinhole. Trajectories of atom’s motion are changed by the action of attractive forces to the walls of nanopinhol channel. In the paraxial approximation the process can be looked upon as an atom beam being defocused by a diverging lens with focal distance:

\[
f_{\text{edW}} \approx -\frac{1}{12} \frac{E_k}{C_3 d} s^5
\]

where \( C_3 \) is the van der Waals coefficient, \( d \) is the thickness of the screen, \( E_k \) is the atom’s kinetic energy. Van der Waals's interaction does not limit the resolution of an atom pinhole camera when the condition \( |f_{\text{edW}}| \gg l \) is fulfilled. This relationship defines the pinhole’s minimum size:

\[
s \gg a_{\text{min}} = \sqrt[5]{\frac{12 C_3 dl}{E_k}}
\]

For example, for atoms of Cs and a silicon screen 50 nm thick the minimum pinhole radius \( a_{\text{min}} \approx 55 \text{ nm} \), for atoms of He the pinhole radius \( a_{\text{min}} \approx 1 \text{ nm} \). The above consideration of atom pinhole camera’s optics shows, that its realization calls for a nanometer diameter pinhole in a screen of nanometer thickness.

In the paper [20] it was demonstrated that nanolithography with use of atom pinhole camera is possible when optimum distance from the pinhole to the mask falls in the range \( L = 1–10 \text{ cm} \), while \( l_{\text{opt}} \approx 10–30 \text{ mkm} \). The “reducing power” of the atom pinhole camera \( M = L / l \) in this case is \( 10^3–10^4 \). For this geometry of atom pinhole camera, typical dimensions of the mask lie within the range of micrometers, and typical dimensions of the structures created on a surface—within the range of nanometers; i.e. atom pinhole camera provides a means for transformation of objects with micrometer sizes into objects with nanometer sizes. There is another outcome of the atom pinhole camera “scaling geometry”, especially important for development of metamaterials. It is the possibility to use in one device not a single pinhole, but their large array. In this case each pinhole generates its own image, which does not intersect the neighboring ones, i.e. the realization of an “atom-multiple pinhole camera” (AMPC) is possible.

AMPC nanolithography opens up wide opportunities for simultaneous generation of great numbers nanostructures for metamaterial fabrication: (1) with nanostructures’ position and size disorder, (2) lithography of identical nanostructures arranged on to the substrate surface in the appropriate ordered way. The first case of AMPC nanolithography was reported in [19]. The second one is the main topic of this paper. Two major questions come to mind in this case: (1) how far identical are the nanostructures; (2) does the distance between neighboring nanostructures (spatial period) is the same on the whole substrate surface. Main limitations to the identity of the parameters in AMPC nanolithography could be divided as attributed to: (1) membrane with nanoholes and (2) surface substrate. Let us note that even at significant quantity of pinholes (up to 10 million ones) inclined-beams aberrations (beginning to show up in the outermost pinholes) are not that restrictive for the resolution of an “atom-multiple pinhole camera”. Thus effect of membrane is restricted to technical limitations of production regularly spaced nanoholes with equal diameters and mechanical stability of thick membrane [22]. Example of the effect attributed to the surface substrate and limiting the identity of fabricated nanostructures is atom’s surface diffusion. It is known that the effect leads to gain lateral size of nanostructures created by surface growth approaches [23]. In the case of AMPC nanolithography it opens a problem of identity of the effective atom-surface sticking coefficient throughout the surface substrate.
3. Experimental setup

Layout drawing of the atom pinhole camera realized in this research is shown in Fig. 1(a). Besides the pinhole itself it includes: an atom beam, a mask, a nanoaperture and a substrate on which the nanostructures were created. The atoms having passed through the mask apertures form, by analogy with ray optics, a “luminous object” of prearranged geometry. Parameters of the atom pinhole camera were chosen for reason of gaining camera’s maximum resolution and a possibility to construct large arrays of surface nanostructures: \( l \approx 20 \text{ mkm}, L \approx 5 \text{ cm}, \) diameters of nanoholes \( d \approx 20 \text{ nm} \). At the parameters chosen an AMPC was operated near a limit of it’s resolving power. In this case the impact of both diffraction of de Broglie atom waves on pinhole parameters and van der Waals forces becomes essential.

To produce an array of nanoholes a dual beam column Quanta 200 3D (FEI company) equipped with ELPHY Quantum (Raith company) electronics and software were used for direct ion beam milling suitable to the problem of nanometrous-range holes (up to molecular size) fabrication in a nanometrous-thin membrane produced in a solid [24]. The method enables to make apertures with diameters down to several nanometers [25]. To produce nanoapertures for atom pinhole camera, 40 nm SiO\(_2\) low stress membranes mounted in the center of a cylindrical disc 3 mm in diameter and 0.2 mm thickness (Ted Pella Inc.) have been used (Fig. 2(a)). Important characteristics of the membrane are superior flatness and stability provided by 200 nm Si\(_3\)N\(_4\) support mesh. The mesh divides 40 nm thick 0.5 mm \( \times \) 0.5 mm SiO\(_2\) film on to 24 fields of 50 \( \mu \text{m} \times 50 \mu \text{m} \) (Fig. 2(b)). The FIB-entry side of the specimen is coated with 10 nm Al in order to prevent charging. Fig. 2(c) shows a SEM image of one of the SiO\(_2\) membrane field with apertures (diameter \( d \approx 80 \text{ nm} \)) arranged in staggered rows.

An AMPC with the above parameters has been realized and employed for fabrication of nanostructures made of In, Au and Ag on a silicon and glass surface. AMPC has been placed into a UHV chamber with the after-vapor pressure in the order of 2 \( \times 10^{-7} \) mbar (Fig. 1(b)). In the experiments a mask was produced from 40 \( \mu \text{m} \) thick metal screen, in which by the method of laser cutting dissimilar-widths through slits was made (Fig. 3(a)). As a source of atomic beam a high temperature effusion cell was used, operated close to the top limit of atomic beam flux applied in the MBE layer growth applications, providing rates of nanostructures growth up to 0.3 Å/s. The generation time for a nanostructures series on one substrate has been determined by atom beam intensity and desired value of nanostructures height. The generation time for a nanostructures series on one substrate has been determined by atom beam intensity and desired value of nanostructures height. Typical time of exposure in the experiment has been \( t \sim 10 \text{ min} \) for nanostructures of height \( h \sim 25 \text{ nm} \). The geometry of nanostructures has been studied by means of atomic force microscope CP-II of the Veeco company.

4. Experimental results

Fig. 3(b) shows AFM image of single nanostructure of In atoms on silicon surface created by the atom pinhole camera with the use of nanoapertures of diameter \( d \approx 20 \text{ nm} \). The presented image shows that form of the nanostructure topologically copies the mask: an individual nanostructure consists of parallel dissimilar-widths stripes crossed by separate stripe, built up from atoms of
Fig. 2. (a) Photo of membrane with holder in the form of a disc 3 mm in diameter and 200 μm in thickness for the atom pinhole camera. (b) SEM image of 40 nm thick membrane divided by 200 nm thick Si3N4 support mesh on to 24 fields. (c) SEM image of one of the SiO2 membrane field with apertures ≈ 70 nm in diameter, manufactured by the method of ion beam milling. In the inset is SEM image of a single nanohole.

In and separated by equal distances of 390 nm. Width of the nanostructure’s first stripe from the left is less than should be in accordance with the width of appropriate slit in the mask; this has been caused by the final aperture of the atom beam, the diameter of which at the mask location is less than the mask size. That is to say, the beam

Fig. 3. Atom nanolithography of a single nanostructure: (a) a photo of the mask used; (b) an AFM image of a nanostructure built up from In atoms on a silicon surface.

atoms have not passed through full aperture formed by the mask’s slits, resulting in curtailment of the stripe length and width. Analysis of the nanostructure’s geometrical parameters Fig. 3(b) indicates that the width and height of its constituent strips differ, being determined by widths of the slits in the utilized mask. To the 250 μm slit there corresponds a nanostructure element with the width of Δ1 ≈ 120 nm and the height of h1 ≈ 7.6 nm, to the 100 μm slit—an element with the width of Δ2 ≈ 80 nm and the height h2 ≈ 2.8 nm. The minimum size of an element, built up in the generated nanostructure from the atoms, that passed through a mask’s slit with the width of 40 μm, equals to 50 nm. This value is 18 nm larger then one obtained in a linear atom trajectories analysis of pinhole camera imagery and attributed to effect of van der
Waals interaction and atom diffraction on to nanoaperture. The height of this element does not exceed the value of 1.8 nm. Non-uniformity of the measured nanostructure’s elements in height is determined by various widths of the slits in the utilized mask, which dictates the stream of atoms building up the corresponding element of the nanostructure.

Possibility of use APC nanolithography for creation of metamaterials has been investigated in a separate experiment. For this purpose nanostructures have been constructed by means of an atom pinhole camera containing an array of nanoapertures with a period 5 μm, see Fig. 4(a). Such an array with holes diameter about 30 nm was created in every 24 fields of the SiO₂ film (Fig. 2(b)). Thus whole area occupied by nanostructures was 24 × 50 μm × 50 μm. Created nanostructures of In atoms on glass surface are presented in Fig. 4(b). This figure shows that the arrangement of nanostructures on the substrate correlates with that of nanoapertures in the membrane of the AMPC: each nanostructure is formed by atoms having passed through a particular nanoaperture.

To explore the identity of nanostructures created we have measured dispersion of stripes width for outermost nanostructures on the substrate. It was found that the value is less than 2% and corresponds to the resolution limit of our AFM operated with ultra sharp 1 nm tip. While dispersion of stripes height was measured on the level of 13% and can be attributed to nonidentity of nanoholes diameter: the height is proportional to intensity of the atom beam near the substructure surface which in turn depends on nanoaperture diameter. These measurements show that at chosen parameters of experiment, when the AMPC is operated near a limit of its resolving power, effect of nanoholes diameter on nanostructures width is negligible in comparison to its impact on to nanostructures height. This is direct evidence of influence both diffraction of de Broglie atom waves on pinhole parameters and van der Waals forces. Measured value for a dispersion of nanostructure’s space period localization on the substrate is appeared to be about 1.6% and is attributed to the thermal drift of the membrane’s holder during ion-beam milling procedure.

Measured difference of nanohole’s diameter is a direct consequence of used technique for fabrication of nanoholes in the membrane. In this method massive ions with energies of thousands electron-volts impinge on a substrate surface and an atomic scale process starts. In this process approximately one atom is removed from the surface for every incident atom thus identity of nanoholes diameter is sensitive to homogeneity of the membrane. Thus the method should be substantially improved to be used in the AMPC nanolithography of identical nanostructures.

AMPC nanolithography opens up opportunities to build bulk metamaterials by creation of 3D nanostructures. It can be implemented by layer by layer deposition and by controlling height of individual element of each nanostructure. Topology of nanostructure’s layers is determined by geometries of an AMPC mask utilized to build these layers.

The AMPC nanolithography has several advantages in comparing with currently used e-beam lithography for metamaterials fabrication. First, AMPC nanolithogra-
phys is a bottom up approach: the desired nanostructures are created directly from required material. Second, the AMPC has a possibility to produce identical nanostructures in massive parallel way, while in direct writing by a focused e-beam every nanostructure is created one by one. Third, a small atomic kinetic energy opens up a possibility to fulfill lithography on to delicate surfaces without its destruction.

One of the AMPC nanolithography features important for metamaterials fabrication is the possibility to create heterostructures. This comes from the fact that atom pinhole camera imagery weakly depends on sort of material used to produce nanostructures. To build heterostructure composed from two layers of materials “A” and “B” it is necessary to use the pinhole camera with double cell MBE source having possibility to evaporate independently material “A” and material “B”. Successive evaporation of these materials through the AMPC’s mask and a pinhole leads to formation of heterostructures. This approach can be extended to production of heterostructures consisting of layers of multiple materials. There are two basic limitations of the approach: the technical one—clogging of the AMPC’s pinhole that reduces number of possible layers of a heterostructures, and the physical one—when AMPC is operated close to it’s ultimate resolution limit difference of physical properties of evaporated materials (mainly the de Broglie wavelength and the atom-surface van der Waals potential) leads to different focal lengths of the AMPC thus planar dimensions of a heterostructure’s layers attributed to different sort of materials become nonidentical.

5. Conclusion

The above results demonstrate a possibility to generate identical nanostructures on a silicon and glass surfaces by means of an AMPC. Forms and sizes of the nanostructures created in this approach are governed by the topology of utilized masks and the size of nanoapertures in the membranes. In the process it is conceivable to control not only planar dimensions of the nanostructure’s elements, but also their height. This circumstance is essential for generation of nanostructures with complexed 3D geometry: the form of nanostructures is defined by the arrangement of apertures forming the mask, while the height of individual elements of nanostructures—by the diameter of these apertures.

We will point out that the method of nanostructures creation presented in this research falls in the category of nanolithographic mask-using ones. In the known methods of lithography, mask elements for nanostructures creation must have dimensions in a nanometrous range and hence their fabrication is a complicated problem both technologically and fundamentally. One advantage of atom pinhole camera utilization for nanolithographic purposes is its feature to generate images with gigantic reduction of the object size—down to 10 thousand times. This makes it possible to use masks of a micrometrous range of dimensions, and their production presents no big problems.

In conclusion, we have successfully implemented the concept of atom pinhole camera as a novel tool for fabrication of metamaterials offering the following merits: (1) it makes possible nanostructures with typical size down to 50 nm; (2) the nanostructures can have an arbitrary prearranged shape; (3) size and form of nanostructures are determined by well-controlled parameters; (4) creation of the great number of identical nanostructures is possible; (5) a variety of materials for nanostructures (atoms, molecules and clusters) is feasible; (6) the method is free from use of a chemically selective etching; (7) in the process of nanostructures creation no destruction of the substrate surface happens.

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